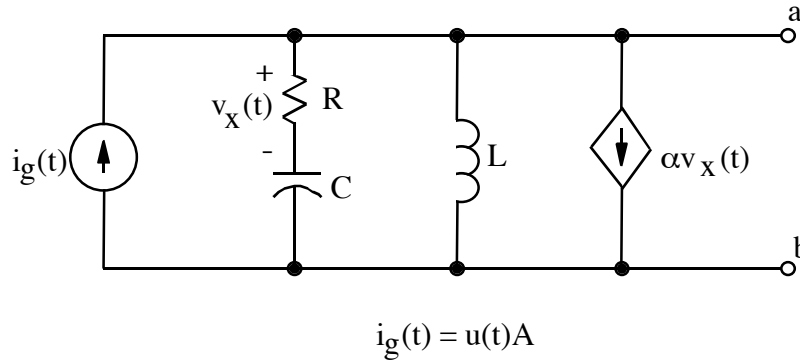
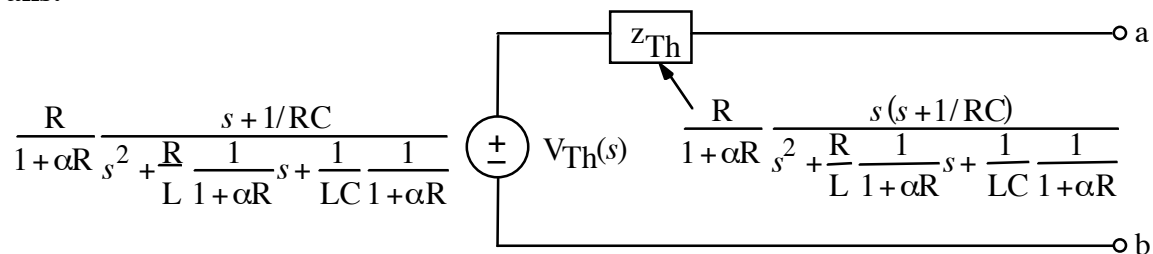


3. (30 points)



Construct an s-domain Thevenin's equivalent to the circuit at the terminals a-b. There is no initial energy stored in the circuit.

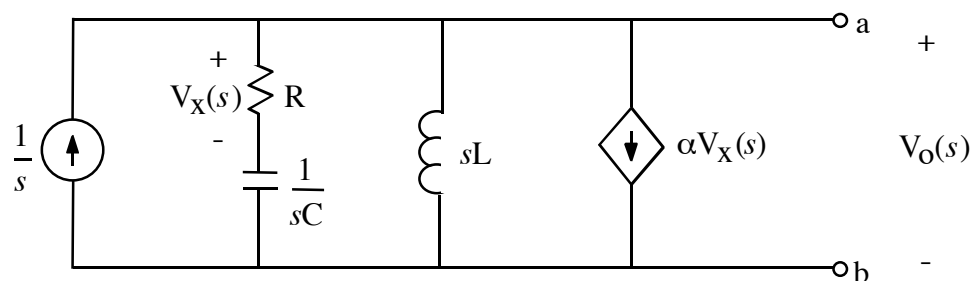
ans:



sol'n: Our s-domain independent current source is

$$\mathcal{L}\{u(t)\} = \frac{1}{s} A$$

No initial energy means we may omit the sources that create initial conditions for L and C in the s-domain. Thus, we have just $1/sC$ and sL :



$V_{Th}(s) = V_O(s)$ with no load across a, b terminals.

We use the node-V method to find $V_o(s)$. The first step is to define the dependent current source in terms of node-voltage $V_o(s)$:

$$\alpha V_x(s) = \alpha V_o(s) \frac{R}{R + \frac{1}{sC}}$$

Aside: Another way to deal with the dependent source is to replace it with an equivalent impedance. From the above equation for $\alpha V_x(s)$, we can say that the dependent source current is equal to $V_o(s)/z_{eq}$ where

$$z_{eq} = \frac{1}{\alpha} \frac{R + \frac{1}{sC}}{R}$$

Having replaced the dependent source with this equivalent impedance, we may then calculate $V_o(s)$ as the independent source current times all the impedances in parallel. Also, we observe that the equivalent impedance is valid when the a,b terminals are shorted out. Thus, this approach is a bit more efficient than the standard node-V method.

For the standard node-V method, we have

$$-\frac{1}{s} + \frac{V_o(s)}{R + \frac{1}{sC}} + \frac{V_o(s)}{sL} + \alpha V_o(s) \frac{R}{R + \frac{1}{sC}} = 0A$$

or

$$V_o(s) \left(\frac{1 + \alpha R}{R + \frac{1}{sC}} + \frac{1}{sL} \right) = \frac{1}{s}.$$

Now we solve for $V_o(s)$ and simplify the result to get a constant times a ratio of polynomials in s with a coefficient of one for the highest power of s in the numerator and denominator. We begin by making each fraction a ratio of polynomials. (We needn't worry yet about the coefficient of the highest power of s .)

$$V_o(s) \left(\frac{s \cdot \frac{1 + \alpha R}{R + \frac{1}{sC}} + \frac{1}{sL}}{s} \right) = \frac{1}{s}$$

$$V_o(s) \left(\frac{s(1 + \alpha R)}{sR + \frac{1}{C}} + \frac{1}{sL} \right) = \frac{1}{s}$$

We put the left side over a common denominator and then move the term in large parentheses to the other side:

$$V_o(s) \left(\frac{s(1 + \alpha R)}{sR + \frac{1}{C}} \cdot \frac{sL}{sL} + \frac{1}{sL} \cdot \frac{sR + \frac{1}{C}}{sR + \frac{1}{C}} \right) = \frac{1}{s}$$

$$V_o(s) = \frac{1}{s} \cdot \frac{sL \left(sR + \frac{1}{C} \right)}{s^2 L(1 + \alpha R) + sR + \frac{1}{C}}$$

Now we cancel a common factor of s top and bottom and pull out a constant to make the coefficient of the highest power of s in numerator and denominator = 1.

$$V_{Th}(s) = V_o(s) = \frac{R}{1 + \alpha R} \frac{s + 1/RC}{s^2 + \frac{R}{L} \frac{1}{1 + \alpha R} s + \frac{1}{LC(1 + \alpha R)}}$$

Check: The units of α are $1/R$ in the original circuit, and that makes αR unitless. ✓

Check: The units of the numerator of the polynomial ratio are $1/\text{sec}$. The units for all terms in the denominator are $1/\text{sec}^2$. The units for the constant out front are Ω , and the units for the independent current source are Amps (which were left out to avoid clutter). Thus, the entire expression has units of $\Omega \text{Asec} = \text{Vsec}$ or V back in the time domain. ✓

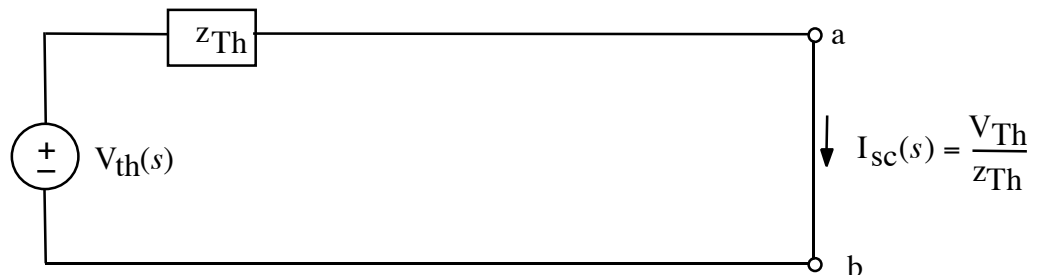
Check: If we set $R = 0$, then we turn off the dependent source and we have

$$V_o(s) = \frac{1}{s} \cdot \frac{1}{sC} \parallel sL = \frac{1}{s} \cdot \frac{L/C}{\frac{1}{sC} + sL} = \frac{1/C}{s^2 + \frac{1}{LC}}.$$

If we plug $R = 0$ into the formula we derived for $V_{Th}(s)$, we get the same answer. (Note that we multiply through by R in the numerator before setting R to zero to avoid a divide by zero.) \checkmark

Check: If we set $L = 0$, then we short out the output terminals and we have $V_o(s) = 0V$. If we plug $L = 0$ into our formula for $V_{Th}(s)$, we get the same answer. (Note that we multiply through by L top and bottom before setting L to zero to avoid a divide by zero.) \checkmark

To find z_{Th} , we use the method of measuring the current, $I_{sc}(s)$, that flows through a short across a, b terminals and setting $z_{Th} = V_{Th}(s)/I_{sc}(s)$:

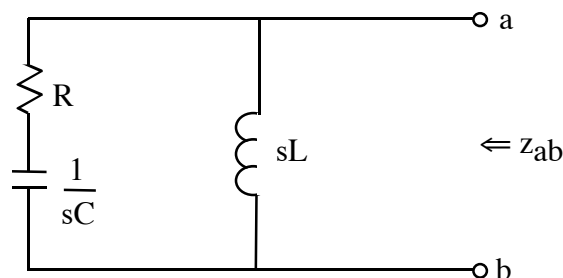


When we perform the same experiment with our actual circuit, all the independent source current flows through the short.

Thus, z_{Th} is just $V_{Th}(s)$ divided by $1/s$. This is the same as multiplying $V_{Th}(s)$ by s :

$$z_{Th} = \frac{R}{1 + \alpha R} \frac{s(s + 1/RC)}{s^2 + \frac{R}{L} \frac{1}{1 + \alpha R} s + \frac{1}{LC} \frac{1}{1 + \alpha R}}$$

Check: If $\alpha = 0$, $z_{Th} = z$ looking into a, b with I-source set to zero. (Dependent source disappears when $\alpha = 0$.)



$$z_{ab} = \frac{s \left(Rs + \frac{1}{C} \right)}{\frac{R}{L}s + \frac{1}{LC} + s^2} = \frac{Rs(s + 1/RC)}{s^2 + \frac{R}{L}s + \frac{1}{LC}}$$

This agrees with our formula for z_{Th} when we plug in $\alpha = 0$. ✓

Check: If $R = 0$, then $V_x(s) = 0$ and dependent source is off. $z_{Th} = z$ looking into a, b with I-source set to zero.

$$z_{ab} = \frac{1}{sC} \parallel sL = \frac{L/C}{\frac{1}{sC} + sL} = \frac{s/C}{s^2 + \frac{1}{LC}}$$

This agrees with our formula for z_{Th} when we plug in $R = 0$, ($R \cdot 1/RC = 1/C$ on top). Also, see earlier check of $V_o(s)$ and multiply by s . ✓

Check: If $C = \infty$, then $1/sC = 0$, $V_x(s) = V_o(s)$, $\alpha V_x(s)$ is the same current we would get with $R_2 = 1/\alpha$.

$$z_{ab} = R \parallel \frac{1}{\alpha} \parallel sL = \frac{R}{1 + \alpha R} \parallel sL = \frac{\frac{R}{1 + \alpha R} \cdot sL}{\frac{R}{1 + \alpha R} + sL}$$

$$z_{ab} = \frac{R}{1 + \alpha R} \frac{s}{s + \frac{R}{L} \frac{1}{1 + \alpha R}}$$

This agrees with our formula for z_{Th} when we plug in $C = \infty$, ($1/\infty = 0$). ✓